

$T > 2.17^\circ\text{K}$ . It was fitted with an empirical formula,  $\beta'(2.200^\circ\text{K}) = 0.30 \times 10^{-3} + (72.0 + 6.66P)^{-1} \text{ atm}^{-1}$ , (3)

to about 1.5%; deviations of the measurements are given in Fig. 3. Also given there is a comparison with  $\beta$  derived from density data of Keesom and Keesom<sup>1</sup> and of Edeskuty and Sherman.<sup>5</sup> Agreement between the three sets of results seems reasonable and the comparison is valid since no  $\lambda$  anomalies exist at this temperature.

A view of Fig. 2 again shows that a compressibility curve between 1.80 and 2.05°K parallels the 2.20°K curve at low pressures, but with increasing pressure it rises above the 2.20°K curve, reaching a peak at  $P_\lambda$ . At  $P > P_\lambda$ , the values of  $\beta$  drop continuously and approach the 2.20°K value. The minima in the  $\beta$ -versus- $P$  curves increase in depth and breadth as  $P_\lambda$  increases, but they seem flattest in the middle of the  $P_\lambda$  range. Although the peaks become more distinct with increased  $P_\lambda$ , the sharpness of all the peak tips required a

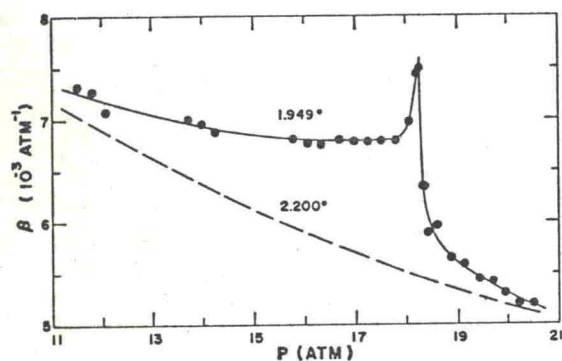


FIG. 4. Compressibility of liquid He<sup>4</sup> at 1.95 and 2.200°K.

higher than normal resolution; therefore the pressure increment of a measurement was reduced from the usual 0.27 to 0.05 atm in the vicinity of the peak. Portions of the 1.95 and 1.80°K curves are shown in Figs. 4 and 5, respectively, along with the 2.20°K values for comparison. The peak at  $P_\lambda$  fades away with increasing temperature until it almost disappears at 2.05°K, although the compressibility excess over the 2.20°K value is still obvious.

In the region between 1.60 and 1.75°K, no  $\lambda$  transition occurs. However, the results in Fig. 6 show that a minimum in the  $\beta$ -versus- $P$  curve persists down to 1.70°K; at 1.75°K, the rate of rise beyond the minimum is similar to that at 1.80°K. Below 1.70°K, (the 1.65°K curve is omitted for clarity) the minimum disappears, but a compressibility excess over the 2.20°K curve remains, amounting to 15% at 1.60°K near the melting pressure.

The temperature variation of  $\beta$  at constant pressure changed according to the proximity of  $(T, P)$  to  $(T_\lambda, P_\lambda)$ . For all temperatures,  $\beta$  at  $P \ll P_\lambda$  increased with in-

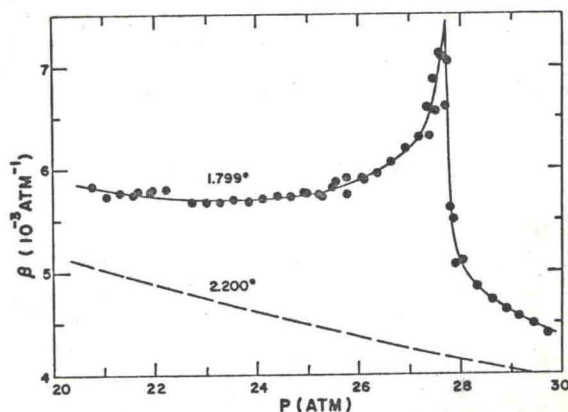


FIG. 5. Compressibility of liquid He<sup>4</sup> at 1.80 and 2.200°K.

creasing temperature. Near the  $\lambda$  transition, the variation of  $\beta$  with temperature became inverted. The reversion of  $(\partial\beta/\partial T)_P$  to the normal plus sign at  $P \gg P_\lambda$  was not indicated—the compressibilities for different temperatures merged to a common value within  $\sim 2\%$ , the experimental error, at the highest pressures.

The accuracy of the measurements is summarized here. From a straight sum of possible individual errors in cell calibration plus those from readings of  $\Delta V$  and  $\Delta P$ , the maximum error in an individual  $\beta$  should be 2.5 to 5.0% for high to low values of  $\beta$ , respectively; from the root mean square of individual errors, a probable error in  $\beta$  is 1.5 to 3.0% for high to low values. Consideration of  $\Delta P_L/\Delta P_U$  alone in Eq. (2) leads to a precision error of 1.0 to 1.7% for high to low values of  $\beta$ . Near the  $\lambda$  transition, the decrease in  $\Delta P$  for greater resolution probably lowered the accuracy, but here we are mainly interested in the reproducibility of results over a short range of pressure and time. Error in these results is estimated at 2%.

#### IV. DISCUSSION

The present compressibility measurements provide a view of normal and abnormal behavior in liquid He<sup>4</sup>

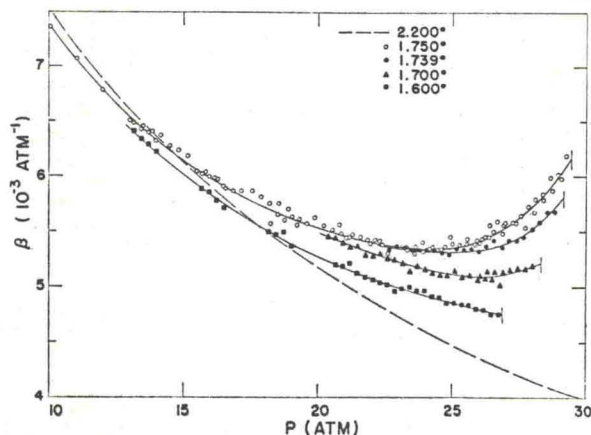
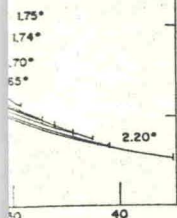


FIG. 6. Compressibility of liquid He<sup>4</sup> at several temperatures below 1.76 and at 2.200°K.

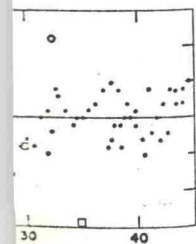


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S.) 64A, No. 1 (1960).